

An Optimization on Quadrature Formulas and Numerical Solutions of Ordinary Differential Equations

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Abstract—This paper is a continuation of the discussion on optimization of the quadrature formulas and their applications in paper [2]. Second-order numerical solutions of Volterra integral equations are constructed using the quadrature formulas obtained in [2]. The numerical results presented in the paper confirm the effectiveness of the methods for numerical solution of ordinary differential equations.

I. INTRODUCTION

IN PAPER [2] we study the quadrature formulas which have generating functions $G_1(x) = \pi \sec(\pi\sqrt{x}/2)/4$ and $G_2(x) = \pi \tan(\pi\sqrt{x}/2)/(4\sqrt{x})$. We construct the second-order quadrature formulas

$$\frac{h}{2} \left(y_0 + \sum_{k=1}^{N-1} \bar{E}_k y_{N-k} + \frac{\pi-1}{2} y_N \right) = \int_a^b y(x) dx + O(h^2),$$

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$$\frac{h}{2} \left(y_0 + \sum_{k=1}^{N-1} \bar{B}_k y_{N-k} + \frac{\pi^2-6}{4} y_N \right) = \int_a^b y(x) dx + O(h^2),$$

where $h = (b-a)/N$ and $x_k = a + kh$, $y_k = y(x_k)$ for $k = 0, 1, \dots, N$ and \bar{E}_k and \bar{B}_k are the coefficients of the Maclaurin series of the generating functions

$$\bar{E}_k = \frac{|E_{2k}|}{(2k)!} \left(\frac{\pi}{2} \right)^{2k+1}, \bar{B}_k = \frac{(4^{k+1}-1)\pi^{2k+2}|B_{2k+2}|}{(2k+2)!}, \quad (1)$$

where E_k and B_k are the Euler and Bernoulli numbers. The coefficients of the right endpoint expansion formulas are equal to the coefficients of the Maclaurin series of the functions $H_i(x) = G_i(e^{-x})$ for $i = 1, 2$. In [2] we construct third-order and fourth-order quadrature formulas as linear combinations of the trapezoidal rule

$$\frac{h}{2} \left(y_0 + 2 \sum_{k=1}^{N-1} y_{N-k} + y_N \right) = \int_a^b y(x) dx + O(h^2), \quad (2)$$

the two quadrature formulas above and their modifications. The method for construction of quadrature formulas by first specifying the generating function is an effective method for construction of approximations of the fractional derivatives and integrals (see [1], [3], [4], [5]). The method is applicable

for construction of approximations of the definite integral and the integer order derivatives of a function as well. In [7] we construct approximations of the first derivative which have exponential and logarithmic generating functions. In the paper we give a proof for the convergence of the approximations and applications of the approximations for numerical solution of ordinary and partial differential equations. The method used in the paper can be extended for construction of approximations of the second derivative which are suitable for deriving approximations of the fractional derivatives. Other methods for numerical solution of integral equations use Gaussian quadratures on non-uniform nets and Monte Carlo methods for numerical integration (see. [8], [9], [10]). Let

$$\frac{h}{2} \sum_{i=0}^N w_i y_{N-i} = \int_a^b y(t) dt + O(h^2) \quad (*)$$

be a second-order quadrature. Consider a Volterra integral equation of the second kind which has the following form

$$y(x) + \int_0^x K(x-t)y(t) dt = F(x). \quad (3)$$

The Nyström method (see [6]) for computing the numerical solution of equation (3) uses the approximations of the values of the definite integral in equation (3) with quadrature formula (*) on all intervals $[0, x_n]$,

$$y_n + \frac{h}{2} \sum_{i=0}^n w_i K_i y_{n-i} = F_n + O(h^2).$$

The numerical solution $\{u_n\}_{n=0}^N$ of integral equation (3), where u_n is an approximation of the value of the solution y_n , is computed recursively with

$$u_n = \frac{1}{2 + w_0 K_0 h} \left(2F_n - h \sum_{i=1}^n w_i K_i y_{n-i} \right) \quad (\text{NS1}^*)$$

and has an initial condition $u_0 = y_0 = F(0)$. The computation of numerical solution NS(*) involves $O(N^2)$ operations. Denote by (3.1) and (3.2) the equations obtained from equation (3) with a kernel function $K(s) = 3 + 2s$ and right hand side $F_1(x) = x^3(10(4x^2 + 30x + 40) \ln x - 18x^2 - 75x)/400$ and

$$F_2(x) = \arctan x - x + (3/2 + x) \ln(1 + x^2) + (1 + 3x + x^2) \operatorname{arccot} x$$

respectively. Equations (3.1) and (3.2) have the solutions $y(x) = x^3 \ln x$ and $y(x) = \operatorname{arccot} x$. In table I we give the experimental results for the error and the order of the numerical solution NS(2), which uses the trapezoidal rule (2), of equations (3.1)-left and (3.2)-right on the interval $[0, 1]$. The rest of the paper is organized as follows. In section two we construct the numerical solution of first order ODEs by transforming them to integral equations in the form (3) and we also construct the numerical solutions which use the corresponding shifted quadrature formulas. In section three of the paper we apply the method for numerical solution of second order ordinary differential equations (ODEs) by converting them to Volterra integral equations

TABLE I

h	$Error$	$Order$	$Error$	$Order$
0.005	4.545×10^{-6}	2.000	2.802×10^{-6}	2.000
0.0025	1.136×10^{-6}	2.000	7.006×10^{-7}	2.000
0.00125	2.841×10^{-7}	2.000	1.751×10^{-7}	2.000

II. NUMERICAL SOLUTION OF FIRST ORDER ODES

In this section we construct the numerical solution of first order ODEs by first transforming them to integral equations which are solved with Nyström method. In [2] we obtain the second order approximations

$$\frac{h}{2} \sum_{k=0}^N \bar{E}_k y_{N-k} = \int_a^b y(x) dx + O(h^2), \quad (4)$$

$$\frac{h}{2} \sum_{k=0}^N \bar{B}_k y_{N-k} = \int_a^b y(x) dx + O(h^2), \quad (5)$$

where $\bar{E}_0 = (\pi - 1)/2$, $\bar{B}_0 = (\pi^2 - 6)/4$ and the rest of the weights are defined with (1). The two approximations (4) and (5) require that the integrand function satisfies $y(a) = 0$. We use the method (see [3], [7]) for extending the approximations to the class of all differentiable functions by changing the values of the last weights. By applying approximation (4) to the function $y(x) - y(a)$ we obtain

$$\begin{aligned} \frac{h}{2} \sum_{k=0}^{N-1} \bar{E}_k (y_{N-k} - y_0) &= \int_a^b (y(x) - y(a)) dx + O(h^2), \\ \frac{h}{2} \left(\sum_{k=0}^{N-1} \bar{E}_k y_{N-k} + 2N y_0 - y_0 \sum_{k=0}^{N-1} \bar{E}_k \right) &= \int_a^b y(x) dx + O(h^2). \end{aligned}$$

Define $\bar{E}_N = 2N - \sum_{k=0}^{N-1} \bar{E}_k$. Therefore approximation (4) has a second-order accuracy for all differentiable functions. Similarly, when $\bar{B}_N = 2N - \sum_{k=0}^{N-1} \bar{B}_k$ approximation (5) holds for all differentiable functions in $[a, b]$. Now we apply approximations (4) and (5) for numerical solution of ODEs. Consider the first order linear ODE

$$y' + ay = f(x), \quad y(0) = y_0. \quad (6)$$

By applying integration on both sides of equation (6) we get

$$\begin{aligned} \int_0^x y'(t) dt + a \int_0^x y(t) dt &= \int_0^x f(t) dt, \\ y(x) + a \int_0^x y(t) dt &= F(x), \end{aligned} \quad (7)$$

where $F(x) = y_0 + \int_0^x y(t) dt$. Equation (7) is equivalent to (6) and it is a Volterra integral equation of the second kind with a kernel $K(s) = a$. Denote by (7.1) the equation which corresponds to (7) when $a = 2$ and $F_1(x) = x^3(4(2 + x) \ln x - x)/8$. Equation (7.1) has the solution $y(x) = x^3 \ln x$ and it is equivalent to equation (6) with a right hand side $f_1(x) = x^2(1 + (3 + 2x) \ln x)$. Denote by (7.2) the integral equation which corresponds to equation (7) with $a = 3$,

$F_2(x) = (1+3x) \operatorname{arccot} x + 3 \ln(1+x^2)/2$ and initial condition $y(0) = y_0 = F_0(0) = \pi/2$. Equation (7.2) has the solution $y(x) = \arccos x$ and it is equivalent to equation (6) with a right hand side $f_2(x) = 3 \operatorname{arccot} x - 1/(1+x^2)$. Equations (7.1) and (7.2) are integral equations in the form (3) and can be solved numerically with methods NS(4) and NS(5), which use quadrature formulas (4) and (5). In table II we give the experimental results for the error and the order of the numerical solution NS(4) of equations (7.2) and (7.3).

TABLE II

h	<i>Error</i>	<i>Order</i>	<i>Error</i>	<i>Order</i>
0.005	7.202×10^{-6}	1.995	2.918×10^{-6}	1.969
0.0025	1.804×10^{-6}	1.997	7.412×10^{-7}	1.977
0.00125	4.513×10^{-7}	1.999	1.870×10^{-7}	1.987

The results for the error and order of the numerical solution NS(5) of equations (7.2) and (7.3) are given in table III.

TABLE III

h	<i>Error</i>	<i>Order</i>	<i>Error</i>	<i>Order</i>
0.005	3.823×10^{-6}	2.004	1.351×10^{-6}	1.996
0.0025	9.544×10^{-7}	2.002	3.385×10^{-7}	1.997
0.00125	2.384×10^{-7}	2.001	8.476×10^{-8}	1.998

Now we construct the shifted quadratures which correspond to formulas (4) and (5). From [2]

$$\frac{h}{2} \sum_{k=0}^N \bar{E}_k y_{N-k} = \int_a^b y(x) dx + \frac{1}{4} f(b)h + O(h^2), \quad (8)$$

$$\frac{h}{2} \sum_{k=0}^N \bar{B}_k y_{N-k} = \int_a^b y(x) dx + \frac{3}{4} f(b)h + O(h^2), \quad (9)$$

where the weights \bar{E}_k and \bar{B}_k are defined with (1) for all indices $k = 0, 1, \dots, n$. From the mean value theorem we have the second order approximation

$$\int_a^{b+ch} y(x) dx = \int_a^b y(x) dx + cf(b)h + O(h^2). \quad (10)$$

From (8) and (10)

$$\frac{h}{2} \sum_{k=0}^N \bar{E}_k y_{N-k} = \int_a^{b+h/4} y(x) dx + O(h^2). \quad (11)$$

Shifted quadrature (11) has a requirement that the integrand function satisfies the condition $y(a) = 0$. By applying (11) to the function $y(x) - y(a)$ we find

$$\frac{h}{2} \sum_{k=0}^{N-1} \bar{E}_k (y_{N-k} - y_0) = \int_a^{b+h/4} (y(x) - y(a)) dx + O(h^2),$$

Shifted quadrature (11) has a second order accuracy for all differentiable functions when the weight \bar{E}_N is defined as

$$\bar{E}_N = \frac{2}{h} \left(b - a + \frac{h}{4} \right) - \sum_{k=0}^{N-1} \bar{E}_k = 2N + \frac{1}{2} - \sum_{k=0}^{N-1} \bar{E}_k$$

Similarly from (9) and (10) we obtain

$$\frac{h}{2} \sum_{k=0}^N \bar{B}_k y_{N-k} = \int_a^{b+3h/4} y(x) dx + O(h^2). \quad (12)$$

Shifted quadrature (12) has a second-order accuracy when

$$\bar{B}_N = 2N + \frac{3}{2} - \sum_{k=0}^{N-1} \bar{B}_k$$

and the weights $\bar{B}_0, \bar{B}_1, \dots, \bar{B}_{N-1}$ are defined with (1). The first weights $\bar{E}_0 = \pi/2$ and $\bar{B}_0 = \pi^2/4$. Now we construct the numerical solution of integral equation (7) which uses shifted quadrature (11). By approximating the definite integral in (7) at the point $x_{n+1/4}$ with (11) we obtain

$$y_{n+1/4} + \frac{1}{2} ah \sum_{k=0}^{n-1} \bar{E}_k y_{n-k} = F_{n+1/4} + O(h^2).$$

Let $\{u_n\}_{n=0}^N$ be the numerical solution of (7). From the second order approximation

$$y_{n+1/4} = \frac{5y_n - y_{n-1}}{4} + O(h^2)$$

we obtain the recursive formula for the numerical solution

$$u_n = \frac{1}{5 + \pi ah} \left(4F_{n+1/4} + u_{n-1} - 2ah \sum_{k=1}^{n-1} \bar{E}_k y_{n-k} \right), \quad (NS1)$$

with an initial condition $u_0 = F_0$. In table IV we give the experimental results for the error and the order of the numerical solution NS1 of equations (7.2)-left and (7.3)-right.

TABLE IV

h	<i>Error</i>	<i>Order</i>	<i>Error</i>	<i>Order</i>
0.005	5.638×10^{-6}	2.000	1.323×10^{-6}	1.996
0.0025	1.409×10^{-6}	2.000	3.314×10^{-7}	1.998
0.00125	3.522×10^{-7}	2.000	8.291×10^{-8}	1.989

The numerical solution of integral equation (7) which uses shifted quadrature (12) is obtained from the approximation of the definite integral in (7) at the point $x_{n+3/4}$ with (12).

$$y_{n+3/4} + \frac{1}{2} ah \sum_{k=0}^{n-1} \bar{B}_k y_{n-k} = F_{n+3/4} + O(h^2).$$

From

$$y_{n+3/4} = \frac{7y_n - 3y_{n-1}}{4} + O(h^2)$$

we obtain the recursive formula of the numerical solution

$$u_n = \frac{2}{14 + a\pi^2 h} \left(4F_{n+3/4} + 3u_{n-1} - 2ah \sum_{k=1}^{n-1} \bar{E}_k y_{n-k} \right). \quad (NS2)$$

The experimental results for the error and the order of the numerical solution NS2 of equations (7.2) and (7.3) are given in table V.

TABLE V

h	$Error$	$Order$	$Error$	$Order$
0.005	6.249×10^{-5}	1.990	1.352×10^{-5}	1.900
0.0025	1.567×10^{-5}	1.995	3.532×10^{-6}	1.936
0.00125	3.924×10^{-6}	1.997	9.054×10^{-7}	1.964

III. NUMERICAL SOLUTION OF SECOND ORDER ODES

We apply the method from section 2 for computing the numerical solution of the second order ODE

$$y'' + 3y' + 2y = f(x), \quad y(0) = y_0, y'(0) = y'_0. \quad (13)$$

Equation (13) is transformed to an integral equation in the form (3) by applying successive integration to both sides

$$\int_0^x y''(u)du + 3 \int_0^x y'(u)du + 2 \int_0^x y(u)du = \int_0^x f(u)du,$$

$$y'(x) + 3y(x) + 2 \int_0^x y(u)du = 3y_0 + y'_0 + \int_0^x f(u)du.$$

Integrate again both sides

$$y(x) + 3 \int_0^x y(t)dt + 2 \int_0^x \int_0^t f(u)dudt = F(x), \quad (14)$$

where $F(x) = y_0 + (3y_0 + y'_0)x + \int_0^x \int_0^t f(u)dudt$. By changing the order of integration of the double integral we get

$$\int_0^x \int_0^t f(u)dudt = \int_0^x \int_u^x f(u)tdtdu = \int_0^x (x-u)f(u)tdtdu.$$

Equation (14) is transformed to

$$y(x) + \int_0^x (3 + 2x - 2t)y(t)dt = F(x). \quad (15)$$

Equation (15) is a Volterra integral equation of the second kind with a kernel $K(s) = 3 + 2s$. In section 1 we compute the numerical solution NS(2) of equations (3.1) and (3.2), which are also integral equations in the form (15). In table VI we give the results for the maximal error and order of numerical solution NS(4), which uses quadrature formula (4), of equations (3.1)-left and (3.2)-right and step sizes $h = 0.005, 0.0025, 0.00125$.

TABLE VI

h	$Error$	$Order$	$Error$	$Order$
0.005	6.730×10^{-6}	1.997	5.885×10^{-6}	1.947
0.0025	1.684×10^{-6}	1.998	1.499×10^{-7}	1.972
0.00125	4.213×10^{-7}	1.999	3.805×10^{-7}	1.978

The results for the error and the order of the numerical solution NS(5) of equations (3.1) and (3.2) are given in table VII.

TABLE VII

h	$Error$	$Order$	$Error$	$Order$
0.005	3.563×10^{-6}	2.002	2.795×10^{-6}	2.008
0.0025	8.901×10^{-7}	2.001	6.969×10^{-7}	2.004
0.00125	2.224×10^{-7}	2.000	1.740×10^{-7}	2.001

IV. CONCLUSION

In the paper we construct second-order numerical solutions of ordinary differential equations by converting them to integral equations and applying the quadrature formulas from [2]. Numerical solutions NS(4) and NS(5) involve additional multiplications compared to the standard method NS(2) and have a longer computational time. One advantage of the methods discussed in the paper is that the linear ODEs are equivalent to Volterra integral equations in the form (3) and they can be solved numerically with NS(*) which uses an appropriate quadrature. All methods discussed in the paper involve $O(N^2)$ computations and have comparable performance. From the experiments presented in the paper and the results of additional experiments we can conclude that the numerical solutions discussed in the paper have a stable and efficient performance. In future work we will prove the convergence of the numerical solutions constructed in the paper and we will apply the methods for numerical solution of other classes of ordinary differential equations.

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